

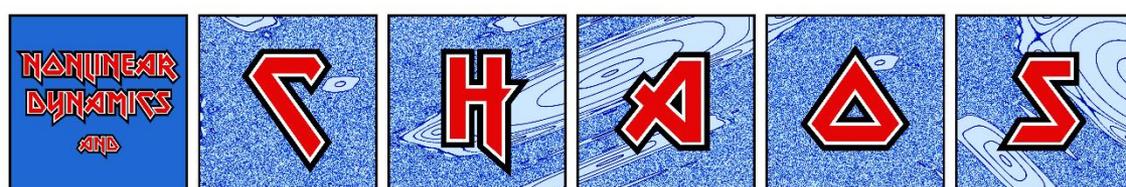
# Lagrangian Descriptors: A powerful method for investigating the behavior and chaoticity of dynamical systems

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# Collaboration

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- **Jérôme Daquin (University of Namur, Belgium)**
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# Outline

- **Lagrangian descriptors (LDs)**
- **Smaller Alignment Index (SALI)**
- **Chaos diagnostics based on LDs:**
  - ✓ **the difference of LDs of neighboring orbits**
  - ✓ **the ratio of LDs of neighboring orbits**
- **Applications:**
  - ✓ **Hénon – Heiles system**
  - ✓ **2D Standard map**
  - ✓ **Motion of a satellite**
- **Summary**

# Lagrangian descriptors (LDs)

The computation of LDs is based on the accumulation of some positive scalar value along the path of individual orbits.

Consider an  $N$  dimensional continuous time dynamical system

$$\dot{\mathbf{x}} = \frac{d\mathbf{x}(t)}{dt} = \mathbf{f}(\mathbf{x}, t)$$

**The arclength definition** [Madrid & Mancho, Chaos (2009) – Mendoza & Mancho, PRL (2010) – Mancho et al., Commun. Nonlin. Sci. Num. Simul. (2013)].

**Forward time LD:**

$$LD^f(\mathbf{x}, \tau) = \int_0^\tau \|\dot{\mathbf{x}}(t)\| dt$$

**Backward time LD:**

$$LD^b(\mathbf{x}, \tau) = \int_{-\tau}^0 \|\dot{\mathbf{x}}(t)\| dt$$

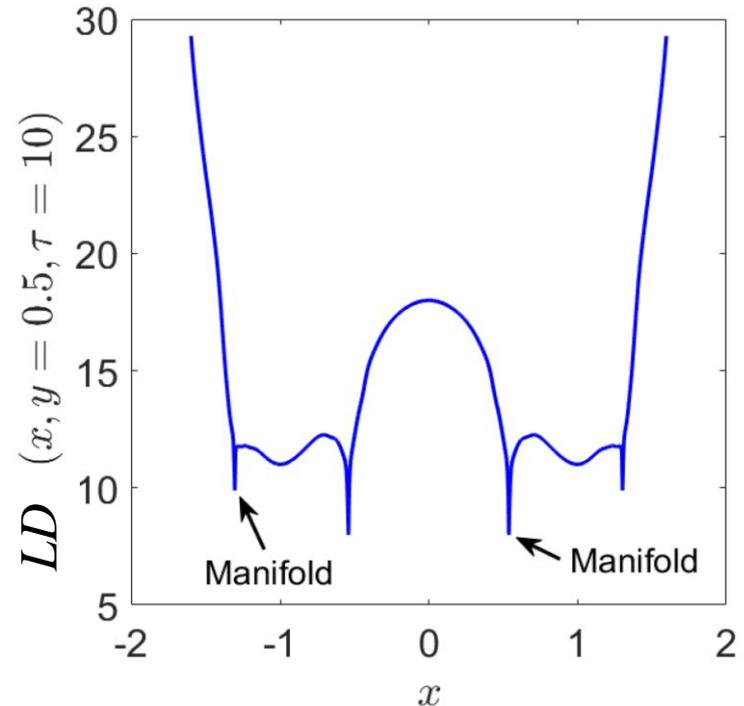
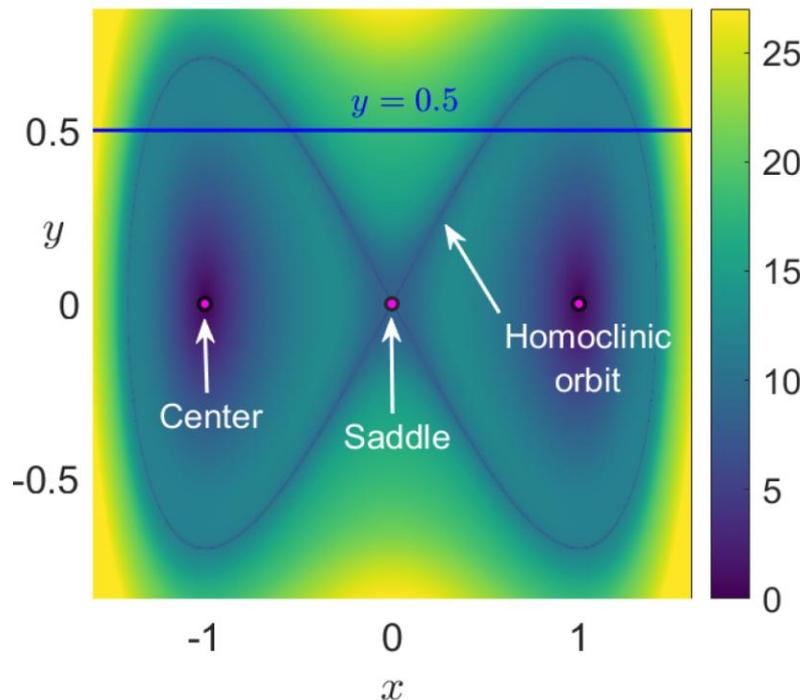
**Combined LD:**

$$LD(\mathbf{x}, \tau) = LD^b(\mathbf{x}, \tau) + LD^f(\mathbf{x}, \tau)$$

# LDs: 1 dof Duffing Oscillator

$$H(x, y) = \frac{1}{2}y^2 + \frac{1}{4}x^4 - \frac{1}{2}x^2$$

The system has three equilibrium points: a saddle located at the origin and two diametrically opposed centers at the points  $(\pm 1, 0)$ .



From Agaoglou et al. 'Lagrangian descriptors: Discovery and quantification of phase space structure and transport', 2020, <https://doi.org/10.5281/zenodo.3958985>

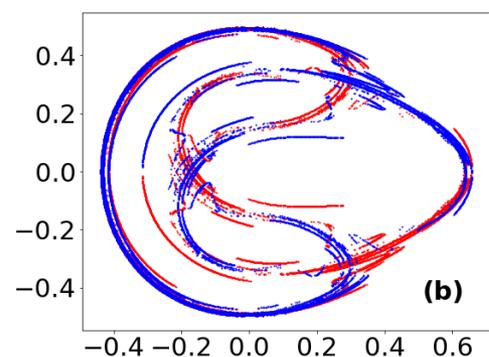
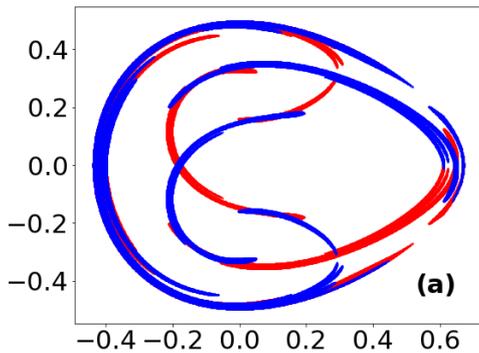
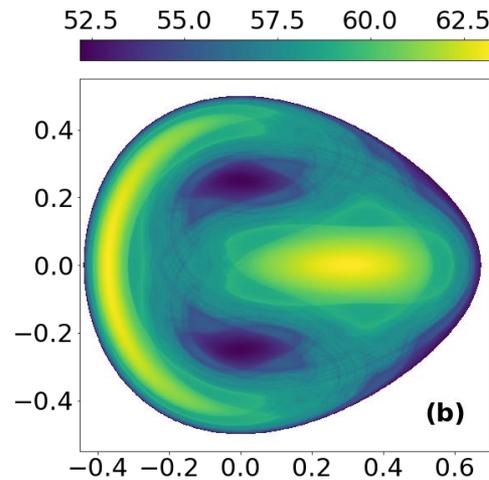
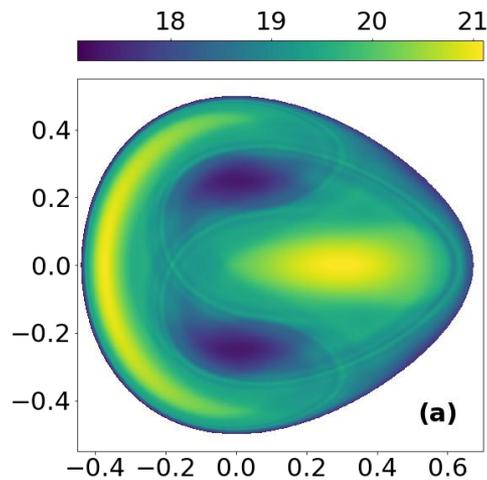
The **location of the stable and unstable manifolds** can be extracted from the ridges of the **gradient field of the LDs** since they are located at points where the forward and the backward components of the LD are non-differentiable.

# Lagrangian descriptors (LDs)

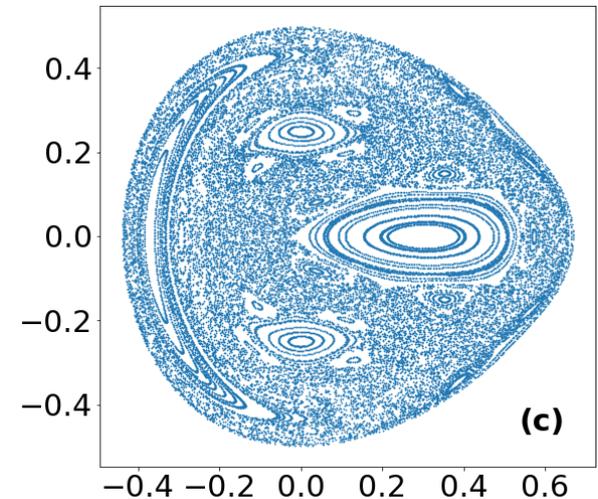
Hénon-Heiles system:  $H = \frac{1}{2}(p_x^2 + p_y^2) + \frac{1}{2}(x^2 + y^2) + x^2y - \frac{1}{3}y^3$

$\tau=20$

$\tau=60$



$H=1/8$



Stable and unstable manifolds

# Lagrangian descriptors (LDs)

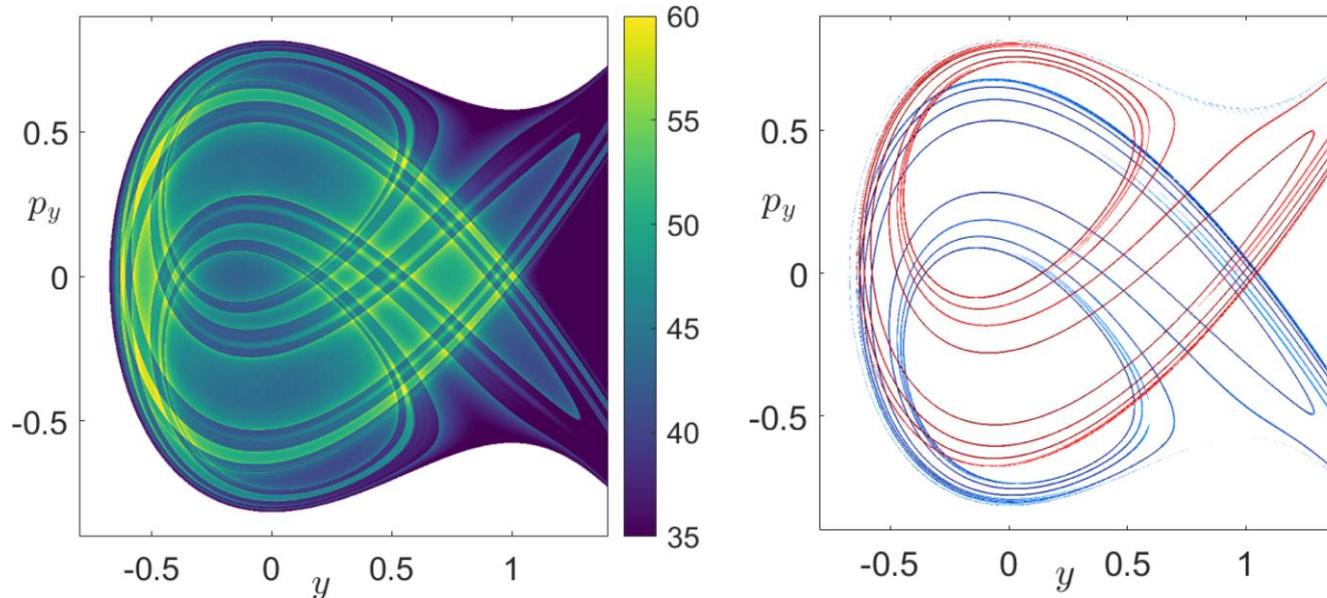
**The ‘ $p$ -norm’ definition** [Lopesino et al., Commun. Nonlin. Sci. Num. Simul. (2015) – Lopesino et al., Int. J. Bifurc. Chaos (2017)].

**Combined LD** (usually  $p=1/2$ ):

$$LD(x, \tau) = \int_{-\tau}^{\tau} \left( \sum_{i=1}^N |f_i(x, t)|^p \right) dt$$

**Hénon-Heiles system:**  $H = \frac{1}{2}(p_x^2 + p_y^2) + \frac{1}{2}(x^2 + y^2) + x^2y - \frac{1}{3}y^3$

**Stable** and **unstable** manifolds for  $H=1/3$ ,  $\tau=10$ .



# Maximum Lyapunov Exponent (MLE)

**Chaos: sensitive dependence on initial conditions.**

Roughly speaking, the MLE of a given orbit characterizes the **mean exponential rate of divergence** of trajectories surrounding it.

Consider an orbit in the  $2N$ -dimensional phase space with **initial condition  $\mathbf{x}(0)$**  and **an initial deviation vector (small perturbation) from it  $\mathbf{v}(0)$** .

Then the mean exponential rate of divergence is:

$$\text{MLE} = \lambda_1 = \lim_{t \rightarrow \infty} \Lambda(t) = \lim_{t \rightarrow \infty} \frac{1}{t} \ln \frac{\|\mathbf{v}(t)\|}{\|\mathbf{v}(0)\|}$$

$\lambda_1 = 0 \rightarrow$  Regular motion ( $\Lambda \propto t^{-1}$ )

$\lambda_1 > 0 \rightarrow$  Chaotic motion

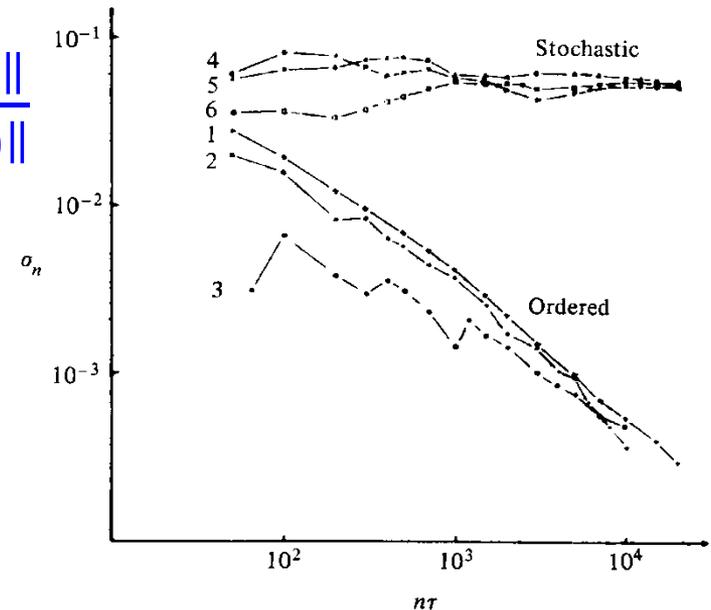


Figure 5.7. Behavior of  $\sigma_n$  at the intermediate energy  $E = 0.125$  for initial points taken in the ordered (curves 1–3) or stochastic (curves 4–6) regions (after Benettin *et al.*, 1976).

# The Smaller Alignment Index (SALI)

Consider the **2N-dimensional phase space** of a conservative dynamical system (**symplectic map or Hamiltonian flow**).

**An orbit** in that space with initial condition :

$$P(\mathbf{0})=(x_1(\mathbf{0}), x_2(\mathbf{0}), \dots, x_{2N}(\mathbf{0}))$$

and a **deviation vector**

$$v(\mathbf{0})=(\delta x_1(\mathbf{0}), \delta x_2(\mathbf{0}), \dots, \delta x_{2N}(\mathbf{0}))$$

The evolution in time (in maps the time is discrete and is equal to the number  $n$  of the iterations) of a **deviation vector** is defined by:

- the **variational equations** (for Hamiltonian flows) and
- the equations of the **tangent map** (for mappings)

# Definition of the SALI

We follow the evolution in time of two different initial deviation vectors ( $v_1(0)$ ,  $v_2(0)$ ), and define SALI [S., J. Phys. A (2001) – S & Manos, Lect. Notes Phys. (2016)] as:

$$\text{SALI}(t) = \min\{\|\hat{v}_1(t) + \hat{v}_2(t)\|, \|\hat{v}_1(t) - \hat{v}_2(t)\|\}$$

where

$$\hat{v}_1(t) = \frac{v_1(t)}{\|v_1(t)\|}$$

When the two vectors become collinear

$$\text{SALI}(t) \rightarrow 0$$

# SALI – Hénon-Heiles system

As an example, we consider the 2D Hénon-Heiles system:

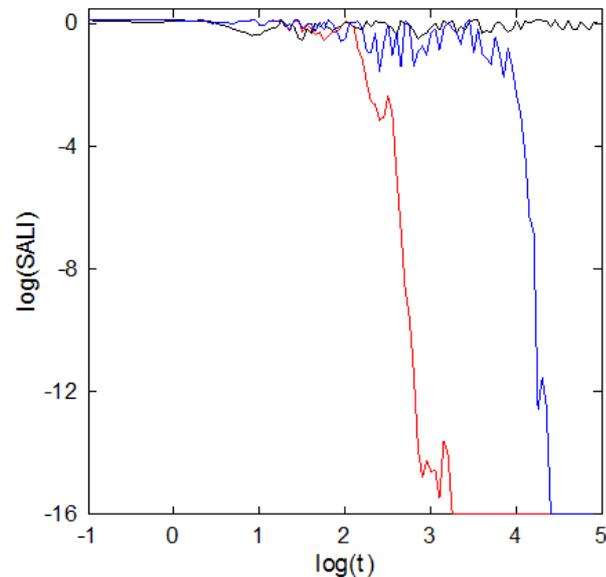
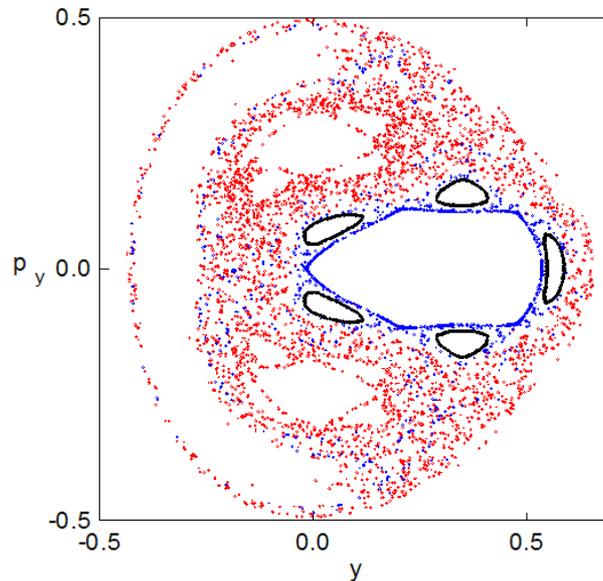
$$H = \frac{1}{2}(p_x^2 + p_y^2) + \frac{1}{2}(x^2 + y^2) + x^2y - \frac{1}{3}y^3$$

For  $E=1/8$  we consider the orbits with initial conditions:

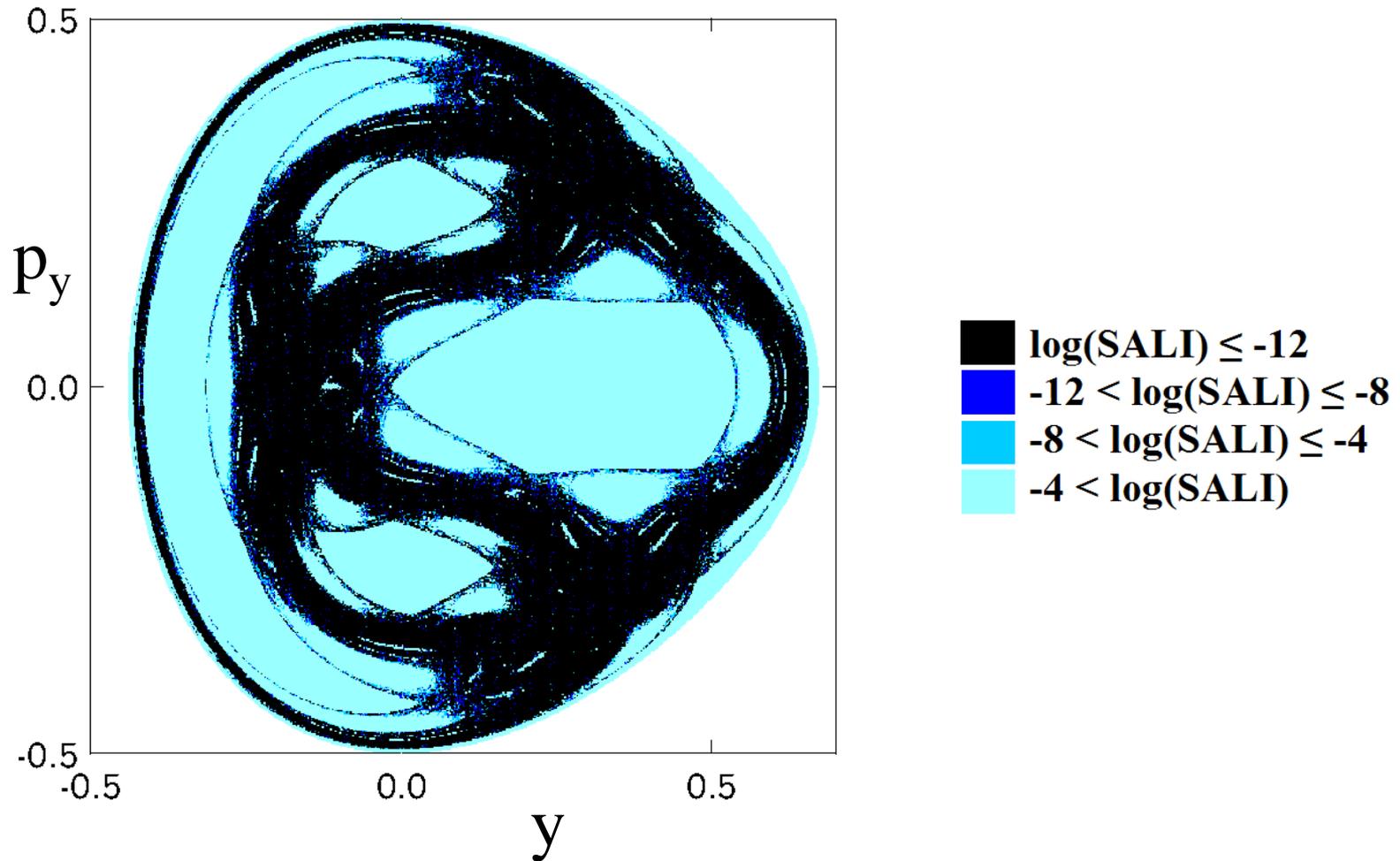
Regular orbit,  $x=0$ ,  $y=0.55$ ,  $p_x=0.2417$ ,  $p_y=0$

Chaotic orbit,  $x=0$ ,  $y=-0.016$ ,  $p_x=0.49974$ ,  $p_y=0$

Chaotic orbit,  $x=0$ ,  $y=-0.01344$ ,  $p_x=0.49982$ ,  $p_y=0$



# SALI – Hénon-Heiles system



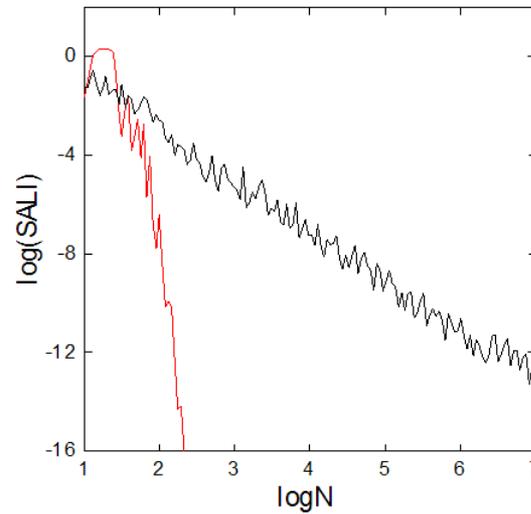
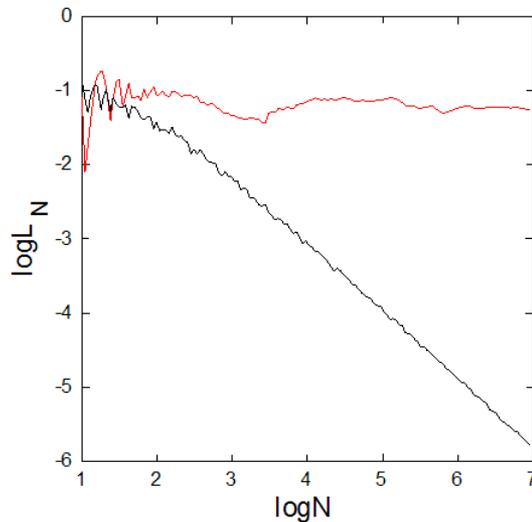
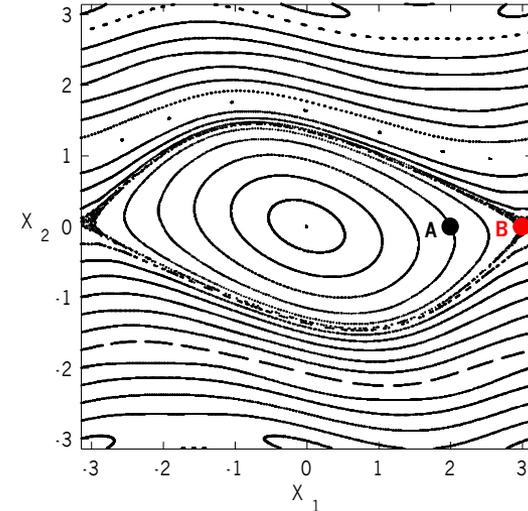
# Applications – 2D map

$$\begin{aligned}x_1' &= x_1 + x_2 \\x_2' &= x_2 - \nu \sin(x_1 + x_2)\end{aligned}\quad (\text{mod } 2\pi)$$

For  $\nu=0.5$  we consider the orbits:

*regular orbit A* with initial conditions  $x_1=2, x_2=0$ .

*chaotic orbit B* with initial conditions  $x_1=3, x_2=0$ .



# Behavior of the SALI

## 2D maps

SALI  $\rightarrow 0$  both for regular and chaotic orbits

following, however, completely different time rates which allows us to distinguish between the two cases.

## Hamiltonian flows and multidimensional maps

SALI  $\rightarrow 0$  for chaotic orbits

SALI  $\rightarrow$  constant  $\neq 0$  for regular orbits

# Using LDs to quantify chaos

We consider orbits on a finite grid of an  $n(\geq 1)$ -dimensional subspace of the  $N(\geq n)$ -dimensional phase space of a dynamical system and their LDs.

Any non-boundary point  $x$  in this subspace has  $2n$  nearest neighbors

$$y_i^\pm = x \pm \sigma^{(i)} e^{(i)}, \quad i = 1, 2, \dots, n,$$

where  $e^{(i)}$  is the  $i$ th usual basis vector in  $\mathbb{R}^n$  and  $\sigma^{(i)}$  is the distance between successive grid points in this direction.

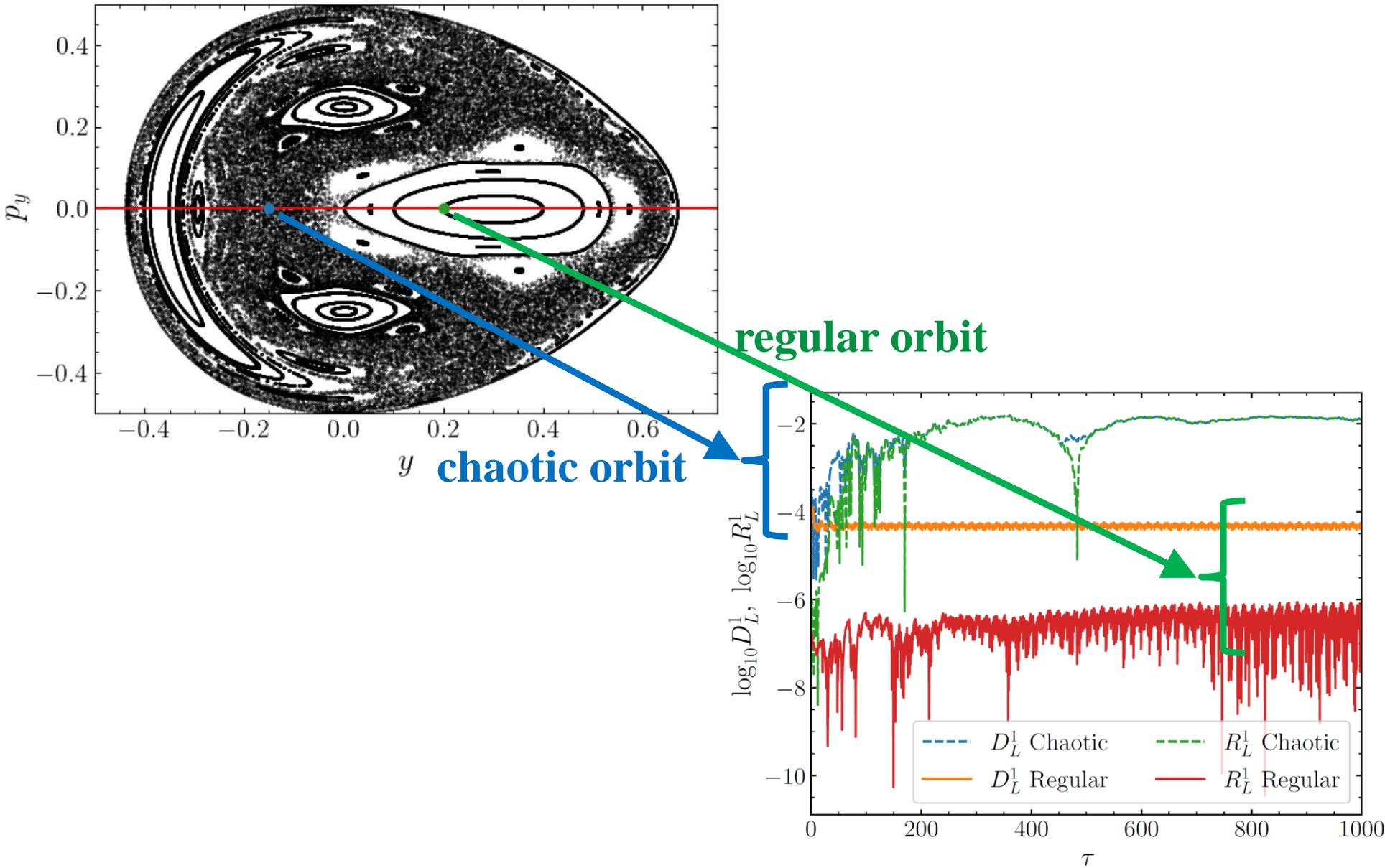
The **difference**  $D_L^n$  of neighboring orbits' LDs:

$$D_L^n(x) = \frac{1}{2n} \sum_{i=1}^n \frac{|LD^f(x) - LD^f(y_i^+)| + |LD^f(x) - LD^f(y_i^-)|}{LD^f(x)}.$$

The **ratio**  $R_L^n$  of neighboring orbits' LDs:

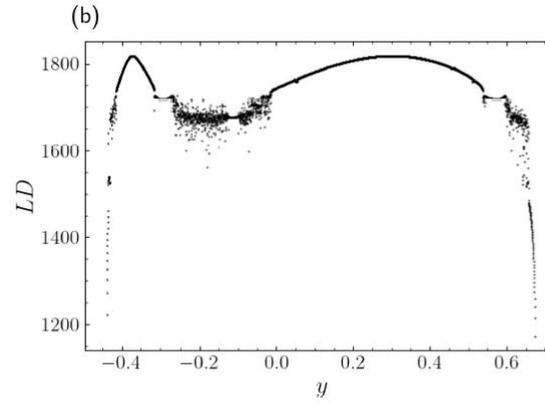
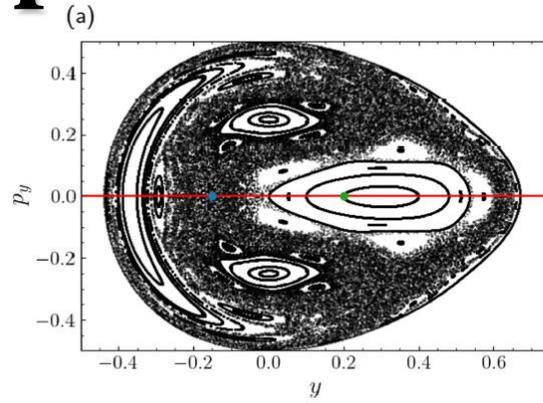
$$R_L^n(x) = \left| 1 - \frac{1}{2n} \sum_{i=1}^n \frac{LD^f(y_i^+) + LD^f(y_i^-)}{LD^f(x)} \right|.$$

# Application: Hénon-Heiles system



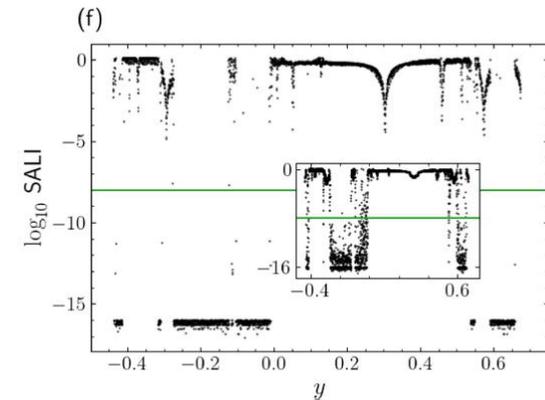
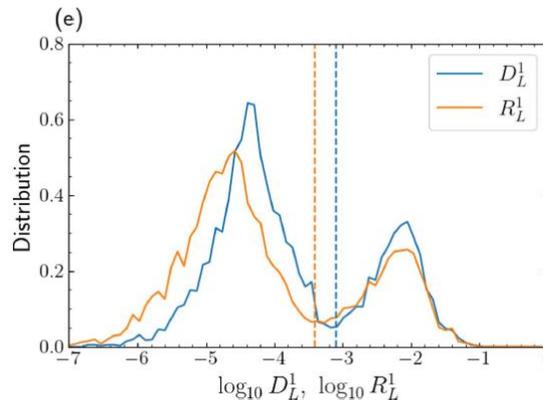
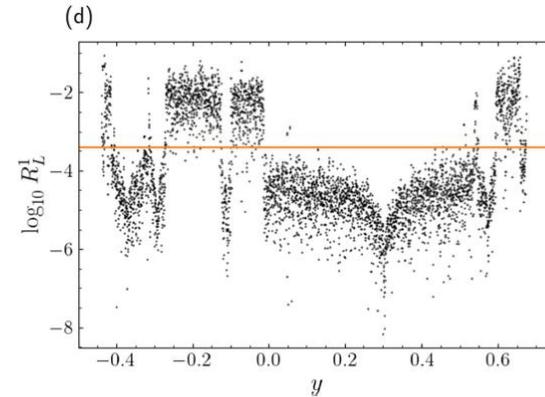
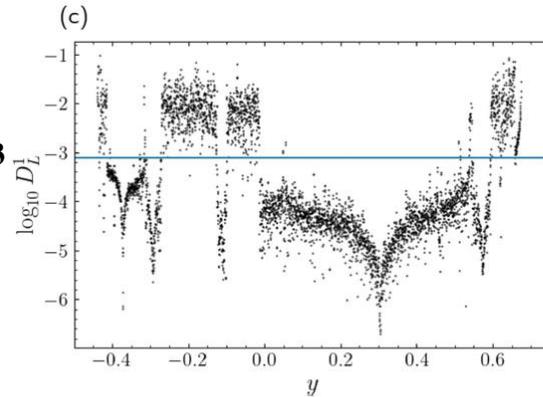
# Application: Hénon-Heiles system

$H=1/8$



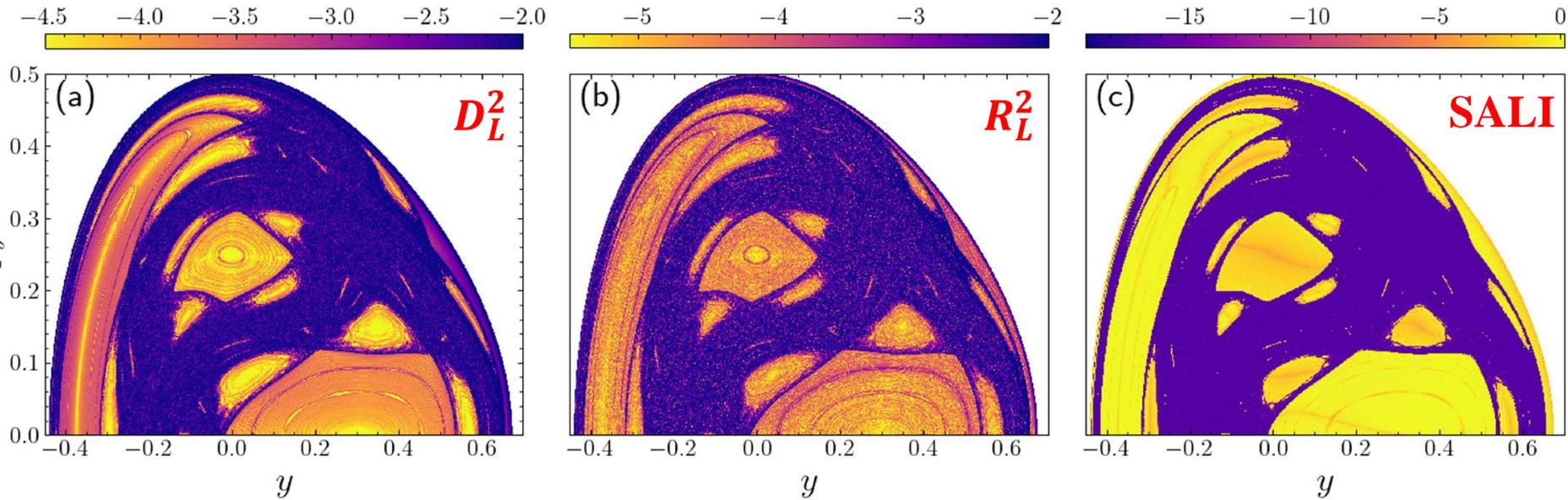
Variation of LDs with regard to initial conditions.  
**regular regions: smooth**  
**chaotic regions: erratic**  
[also see Montes et al., Commun. Nonlin. Sci. Num. Simul. (2021)]

LDs for  $\tau=10^3$

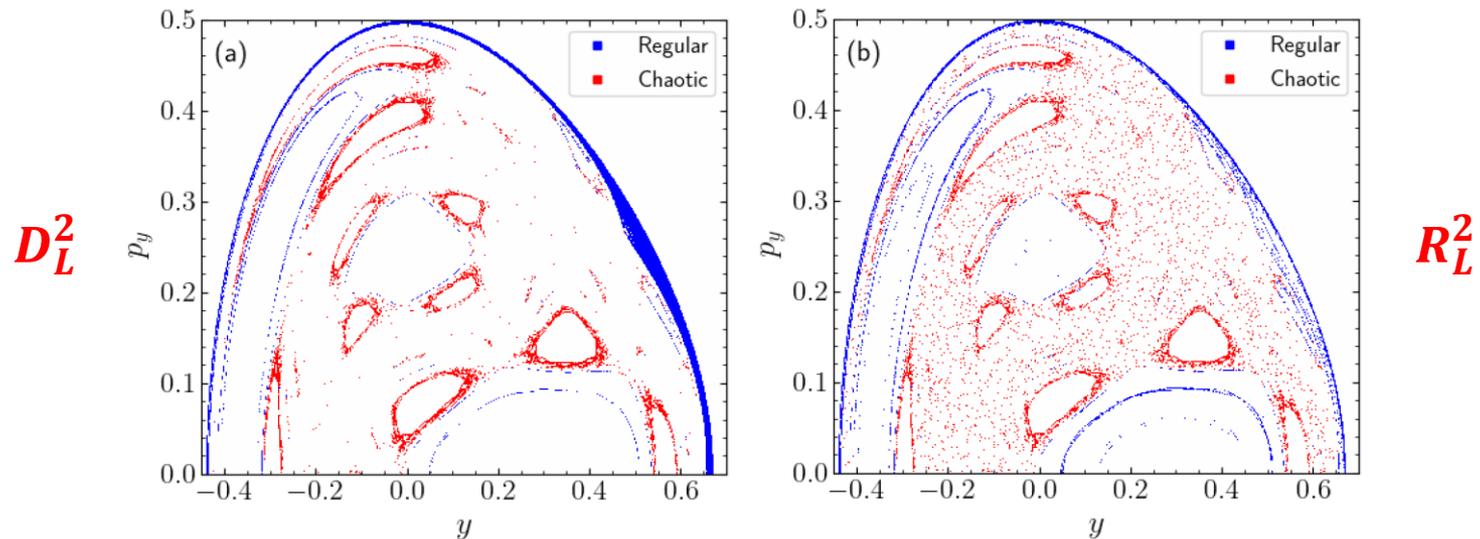


**SALI for  $\tau=10^6$**   
**(inset  $\tau=10^3$ )**

# Application: Hénon-Heiles system



Misclassified orbits (< 10%)

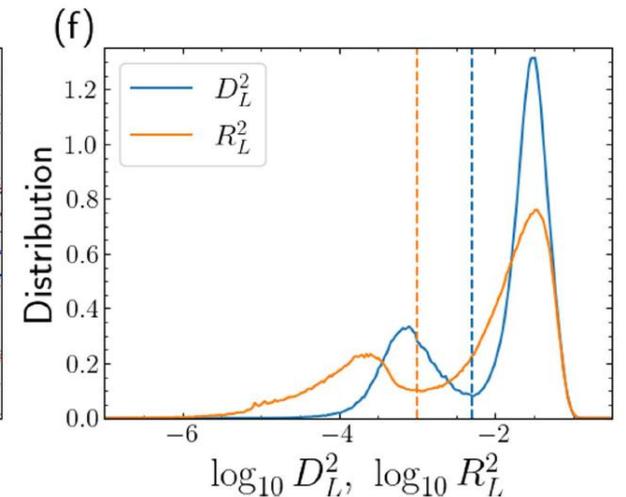
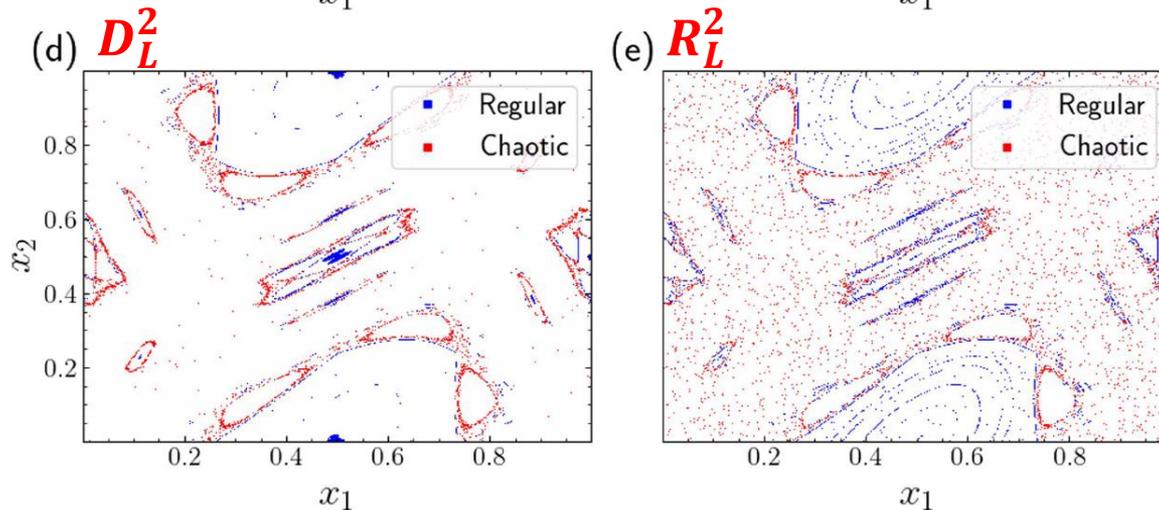
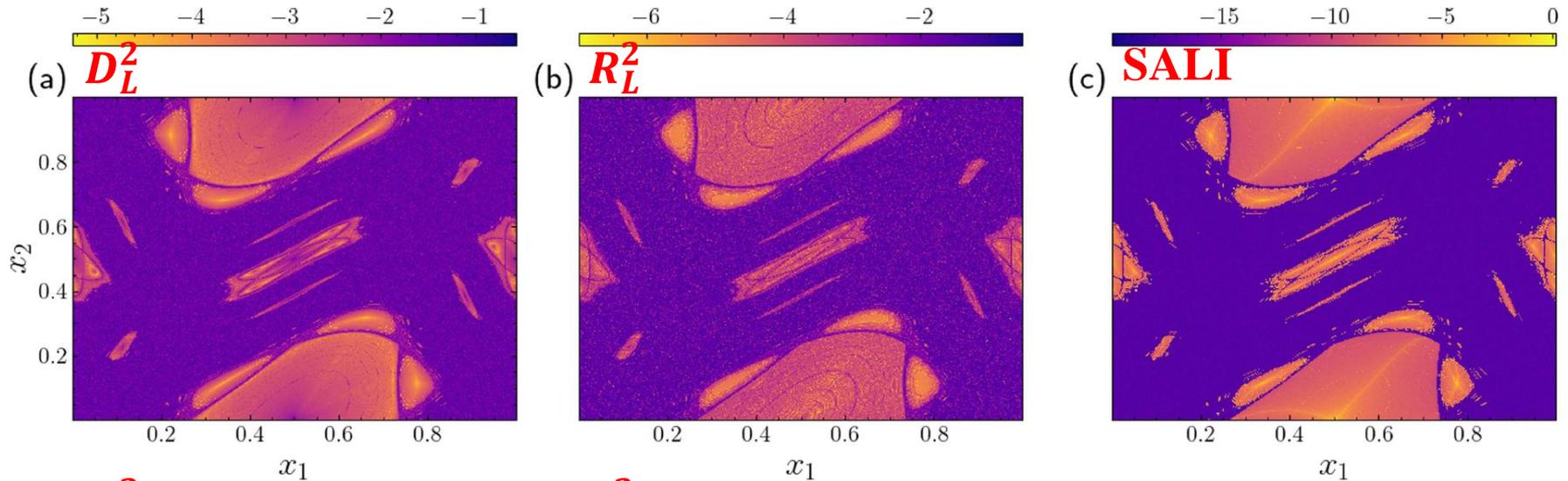


# Application: 2D Standard map

We set  $K = 1.5$

$$\begin{aligned} x'_1 &= x_1 + x'_2 \\ x'_2 &= x_2 + \frac{K}{2\pi} \sin(2\pi x_1) \pmod{1} \end{aligned}$$

Thresholds:  $\log_{10} D_L^2 = -2.3$ ,  $\log_{10} R_L^2 = -3$  ( $T = 10^3$ )  
 $\log_{10} \text{SALI} = -12$  ( $T = 10^5$ )

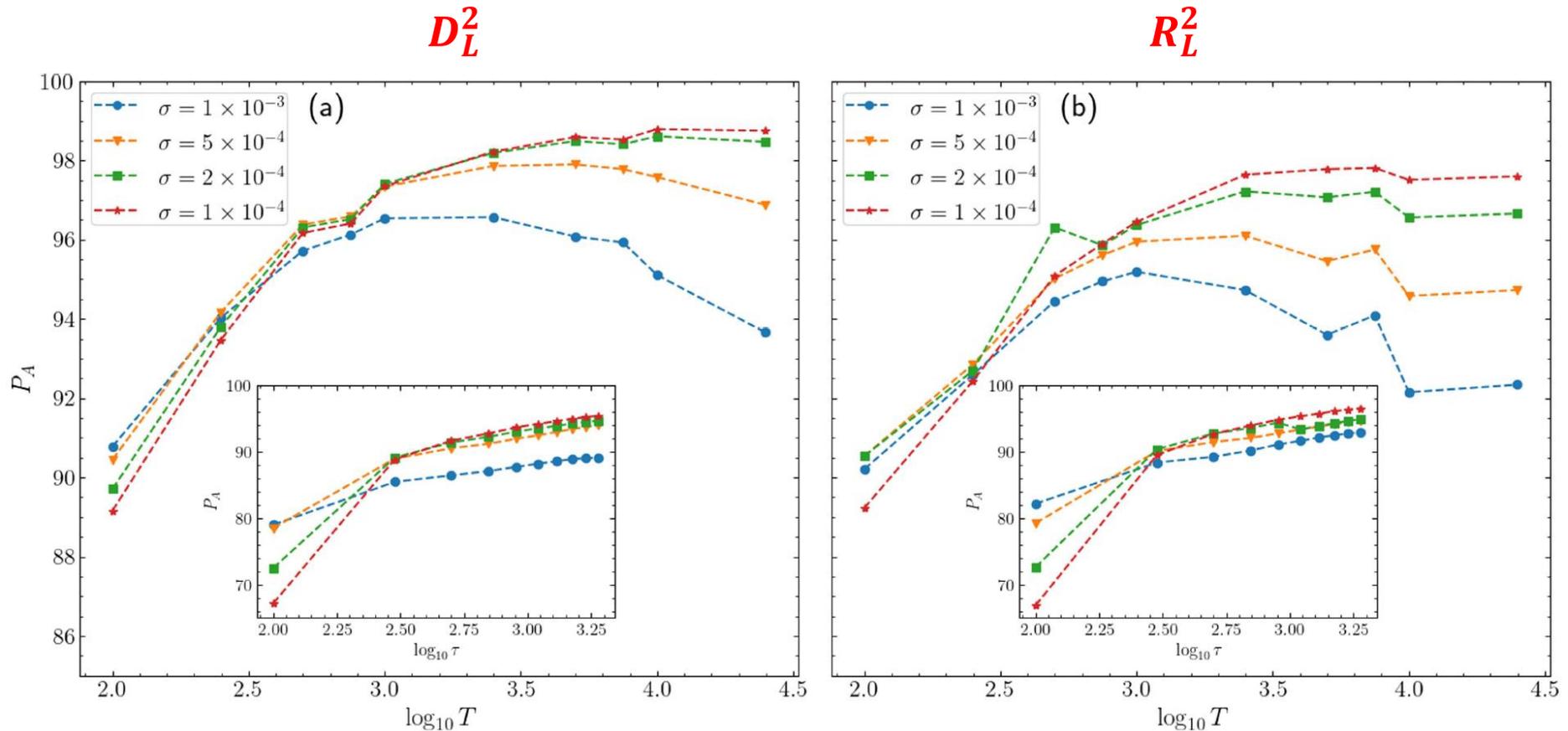


# Effect of grid spacing ( $\sigma$ ) and final integration time ( $T, \tau$ )

$P_A$  : percentage of correctly characterized orbits

Main plots: 2D Standard map

Insets: Hénon-Heiles system

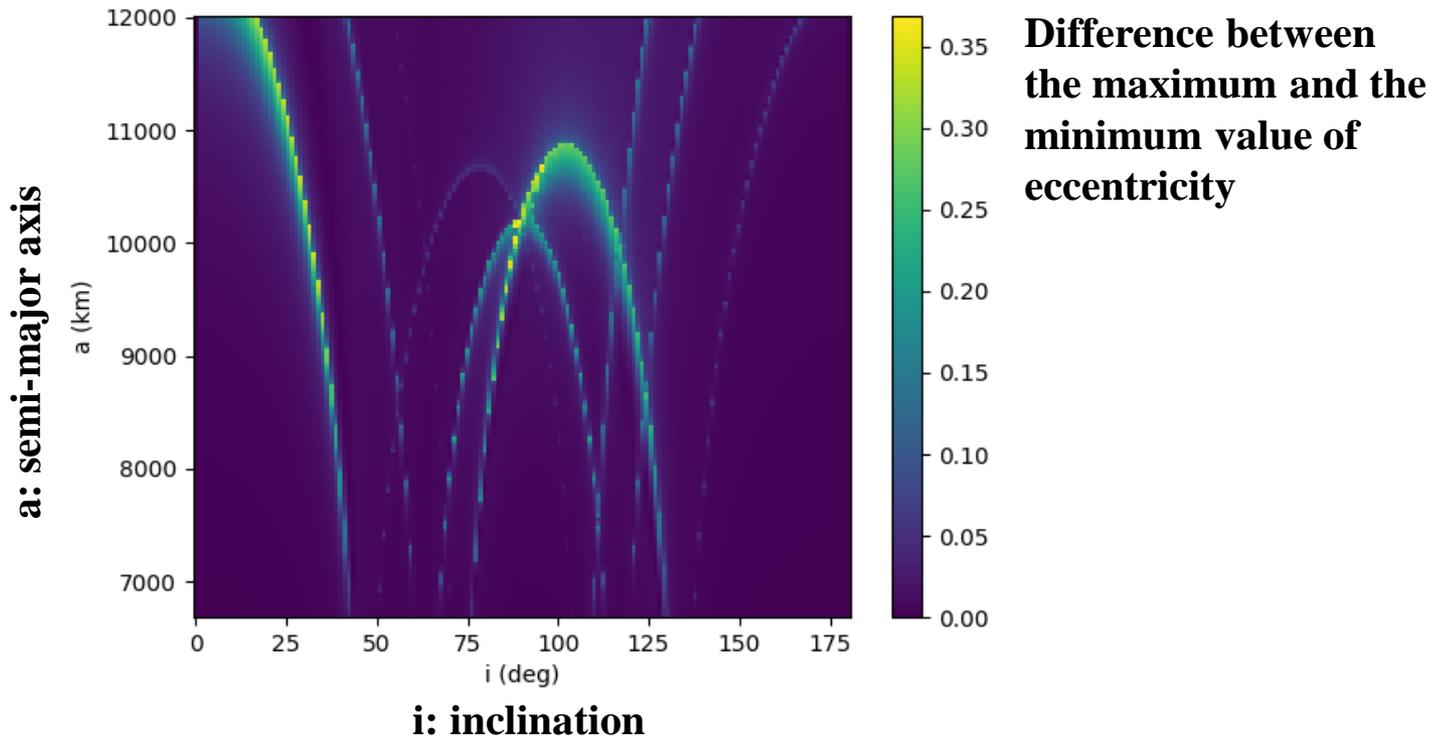


# Current work

Investigation of the **dynamics of a satellite around the Earth.**

Following the work of Gkolias et al., *Cel. Mech. Dyn. Astron.* (2020), we consider a Hamiltonian system consisting of terms describing:

- ✓ **the two body problem** (unperturbed, integrable part), and
- ✓ **perturbation terms due to**
  - **the solar radiation pressure, and**
  - **Earth's oblateness**



# Summary

- ✓ We introduced and successfully implemented computationally efficient ways to **effectively identify chaos** in conservative dynamical systems **from the values of LDs at neighboring initial conditions.**
- ✓ From the distributions of the indices' values we determine appropriate **threshold values**, which allow the characterization of orbits as regular or chaotic.
- ✓ All indices **faced problems** in correctly revealing the nature of some orbits mainly **at the borders of stability islands.**
- ✓ All indices show **overall very good performance**, as their classifications are in accordance with the ones obtained by **the SALI (which is a very efficient and accurate chaos indicator)** at a level of at least **90% agreement.**
- ✓ **Advantages:**
  - **Easy to compute** (actually only the forward LDs are needed).
  - **No need to know and to integrate the variational equations.**

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